An Integrated Numerical Model for the Design of Coastal Protection Structures

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Abstract: In the present work, an integrated coastal engineering numerical model is presented. The model simulates the linear wave propagation, wave-induced circulation, and sediment transport and bed morphology evolution. It consists of three main modules: WAVE_L, WICIR, and SEDTR. The nearshore wave transformation module WAVE_L (WAVE_Linear) is based on the hyperbolic-type mild slope equation and is valid for a compound linear wave field near coastal structures where the waves are subjected to the combined effects of shoaling, refraction, diffraction, reflection (total and partial), and breaking. Radiation stress components (calculated from WAVE_L) drive the depth averaged circulation module WICIR (Wave Induced CIRculation) for the description of the nearshore wave-induced currents. Sediment transport and bed morphology evolution in the nearshore, surf, and swash zone are simulated by the SEDTR (SEDiment TRansport) module. The model is tested against experimental data to study the effect of representative coastal protection structures and is applied to a real case study of a coastal engineering project in North Greece, producing accurate and consistent results for a versatile range of layouts.

Keywords: coastal protection structures; integrated numerical model; waves; hydrodynamics; sediment transport; morphology evolution

1. Introduction

Nowadays, numerical models are the main tool for engineers involved in the design of coastal and marine structures. There are numerous examples in relevant literature of more or less advanced models, covering various aspects of wave-, hydro-, and morpho-dynamics from deep water to the nearshore, at different scales and at varying levels of detail [1–8]. However, coastal engineering practice requires robust integrated models that are able to represent in a reliable way the full range of processes governing coastal dynamics, including the effect of the presence of structures, while maintaining the computational effort needed at reasonable levels. Model versatility should also be considered as an essential requirement, since such models should be able to be adapted to a wide range of design layouts and perform satisfactorily for an equally wide range of field conditions.

For a long time, the design of coastal protection structures was essentially based on engineering experience and empirical rules. Such approaches were gradually replaced by models of varying complexity, focusing on the structures’ effects on wave dynamics, circulation patterns, and morphological evolution in coastal areas. Morphodynamic processes are among the most complex ones to accurately reproduce, since they depend on the combined effect of waves and currents, whose interaction becomes increasingly complicated when moving within the breaker zone and towards the swash. Morphological evolution modelling has indeed come a long way throughout the years, from simple conceptual models to fully 3D ones, currently encompassing a series of improvements in our understanding of the involved processes. Regarding relevant literature (on the more advanced 2D
horizontal, quasi-3D, and 3D models), one may indicatively refer to: pioneering [9] and more recent works [10–12] on bed morphology evolution due to the presence of detached breakwaters, studies on the morphological effects of groins and groin systems [13,14], and more complete/inclusive works on modelling coastal morphodynamics in the presence of coastal structures [15–18]. Focusing on the main features that separate similar integrated modelling attempts, one could refer to: (a) Model capabilities to represent the effects of various types of structures on coastal dynamics; (b) the representation of swash zone dynamics; and (c) the approaches and formulae used to calculate bed load and suspended load sediment transport.

In this work, an integrated coastal engineering numerical model—developed by the authors—is presented and described in detail. The model consists of three main modules that simulate: linear wave propagation (i.e., WAVE_L), wave-induced circulation (i.e., WICIR), and sediment transport along with bed morphology evolution (i.e., SEDTR). The model is capable of simulating the presence of various types of structures (vertical structures; groins and groin systems; emerged, submerged, and floating breakwaters) and includes a novel approach for the representation of swash zone hydrodynamics, while sediment transport is modelled based on the formula proposed by Camenen and Larson [19,20]. The model is tested against experimental data to study the effect of representative coastal protection structures [21,22], and—given the good agreement between calculated results and measured data—is afterwards applied to a real case study of a coastal engineering project in North Greece (combination of submerged breakwaters and beach nourishment).

2. Model Description

2.1. Nearshore Wave Transformation Module–WAVE_L

Linear wave propagation is simulated by applying a mild-slope model [23,24], derived without the assumption of progressive waves. The module WAVE_L is based on the hyperbolic-type mild slope equation and is valid for a compound wave field near coastal structures where waves are subjected to the combined effects of shoaling, refraction, diffraction, reflection (total and partial), and breaking. The module consists of the following pair of equations [23,24]:

\[
\frac{\partial \eta}{\partial t} + \frac{c}{c_g} \nabla \frac{c_g}{c} Q_w = 0
\]

\[
\frac{\partial U_w}{\partial t} + \frac{c^2}{d} \nabla \eta = \nu_h \nabla^2 U_w
\]

where \(\eta\) is the surface elevation; \(U_w\) is the mean velocity vector \(U_w = (U_w, V_w)\); \(d\) is the depth, \(Q_w = U_w h_w = (Q_w, P_w)\); \(h_w\) is the total depth \((h_w = d + \eta)\); \(c\) is the celerity; and \(c_g\) is the group velocity \((c_g = (gd)^{0.5})\). The term \(\nu_h\) is a horizontal eddy viscosity coefficient introduced in order to include breaking effects based on the formulation of [25]:

\[
\nu_h = 2d \left( \frac{D}{\rho} \right)^{1/3}
\]

In Equation (3), \(D\) is the dissipation of wave energy expressed as:

\[
D = \frac{1}{4} Q_b f \rho g H_m^2
\]

where \(H_m\) is the maximum wave height; \(\rho\) is the water density; \(f\) is the wave frequency; and \(Q_b\) is the probability of a wave breaking at a certain depth, expressed as \((1 - Q_b)/(\ln Q_b) = (H_{rms}/H_m)^2\) according to [26]. The mean square wave height \(H_{rms}\) is calculated from \(H_{rms} = 2(\langle 2\eta^2 \rangle)^{1/2}\), with the brackets denoting a time-mean quantity. It should be noted that—since linear wave models are not capable of describing waves in the swash zone—in WAVE_L, the water depth from the rundown point
(i.e., depth equal to $R/4$; $R$ is the runup height) and up to the runup point (i.e., depth equal to $-R$) is considered to be constant and equal to $R/4$.

WAVE-L is adapted for engineering applications based on the following:

1. The input wave is introduced at a line inside the computational domain according to [27,28].
2. A sponge layer boundary condition is used to absorb the outgoing waves at the four sides of the domain [27].
3. The presence of vertical structures is incorporated by introducing a total reflection boundary condition ($U_w = (U_w, V_w) = 0$ normal to the boundary, where $U_w$ is the mean velocity vector; for a rectilinear grid, the above is equivalent to $U_w = 0$ or $V_w = 0$).
4. Partial reflection is also simulated, by introducing an artificial eddy viscosity coefficient $v_h$. The values of $v_h$ are estimated from the method developed by Karambas and Bowers [29], using the reflection coefficient values proposed by Bruun [30].
5. The presence of submerged structures is incorporated as in [31].
6. The presence of floating structures is incorporated as in [32].

The numerical solution is based on the well-documented explicit second order finite difference staggered scheme using a mid-time method [24].

2.2. Wave-Induced Circulation Module—WICIR

The depth and shortwave-averaged 2D continuity and momentum equations are used for simulating nearshore currents in the coastal zone. They are expressed as:

$$\frac{\partial \zeta}{\partial t} + \frac{\partial (U \zeta)}{\partial x} + \frac{\partial (V \zeta)}{\partial y} = 0$$

(5)

$$\frac{\partial U}{\partial t} + U \frac{\partial U}{\partial x} + V \frac{\partial U}{\partial y} + \frac{\partial S_{xy}}{\partial x} = -\frac{1}{\rho h} \left( \frac{\partial S_{xx}}{\partial x} + \frac{\partial S_{xy}}{\partial y} \right) + \frac{1}{h} \frac{\partial}{\partial x} \left( v_h \frac{\partial U}{\partial x} \right) + \frac{1}{h} \frac{\partial}{\partial y} \left( v_h \frac{\partial U}{\partial y} \right) - \frac{\tau_{bx}}{\rho h}$$

(6)

$$\frac{\partial V}{\partial t} + U \frac{\partial V}{\partial x} + V \frac{\partial V}{\partial y} + \frac{\partial S_{xy}}{\partial y} = -\frac{1}{\rho h} \left( \frac{\partial S_{xy}}{\partial x} + \frac{\partial S_{yy}}{\partial y} \right) + \frac{1}{h} \frac{\partial}{\partial x} \left( v_h \frac{\partial V}{\partial x} \right) + \frac{1}{h} \frac{\partial}{\partial y} \left( v_h \frac{\partial V}{\partial y} \right) - \frac{\tau_{by}}{\rho h}$$

(7)

where $S_{xx}$, $S_{xy}$, and $S_{yy}$ are the radiation stresses; $h = d + \zeta$ ($\zeta$ being the mean water elevation); $U$ and $V$ are the depth-averaged current velocities; and $\tau_{bx}$ and $\tau_{by}$ are the bottom shear stresses. Based on linear wave theory, Copeland [33] derived the equations for radiation stresses ($S_{ij}$) without the typical assumption of progressive waves, expressed as:

$$\frac{S_{xx}}{\rho} = d^2 < U_w^2 > A_r - d^2 < \left( \frac{\partial U_w}{\partial x} + \frac{\partial V_w}{\partial y} \right)^2 > B_r + \frac{\partial}{\partial y} \left[ V_w \left( \frac{\partial U_w}{\partial x} + \frac{\partial V_w}{\partial y} \right) \right] > D_r + \frac{1}{2} g \eta^2$$

(8)

$$\frac{S_{yy}}{\rho} = d^2 < V_w^2 > A_r - d^2 < \left( \frac{\partial U_w}{\partial x} + \frac{\partial V_w}{\partial y} \right)^2 > B_r + d^2 \frac{\partial \eta}{\partial y} < \left[ V_w \left( \frac{\partial U_w}{\partial x} + \frac{\partial V_w}{\partial y} \right) \right] > D_r + \frac{1}{2} g \eta^2$$

(9)

$$\frac{S_{xy}}{\rho} = d^2 < U_w V_w > A_r$$

(10)

$$A_r = \frac{k}{4 \sinh^2 kd} (\sinh 2kd + 2kd)$$

(11)

$$B_r = \frac{1}{4 \sinh^2 kd} (\sinh 2kd - 2kd)$$

(12)
\[ D_r = \frac{d}{4 \sinh^2 kd} \left( \frac{1}{2kd} \sinh 2kd - \cosh 2kd \right) \]  

(13)

where \( k \) is the wave number and the brackets denote time-mean quantities.

In nearshore circulation models, the treatment of the bottom stress is critical. The bottom shear stresses \( \tau_b = (\tau_{bx}, \tau_{by}) \) in WICIR are calculated based on the formulae proposed by Kobayashi et al. [34]:

\[ \tau_{bx} = \frac{1}{2} \rho f_b \sigma^2_T G_{bx} \]  

(14)

\[ \tau_{by} = \frac{1}{2} \rho f_b \sigma^2_T G_{by} \]  

(15)

\[ G_{bx} = \frac{U}{\sigma_T} \left[ 1.16^{2} + \left( \frac{|U|}{\sigma_T} \right)^{2} \right]^{0.5} \]  

(16)

\[ G_{by} = \frac{V}{\sigma_T} \left[ 1.16^{2} + \left( \frac{|U|}{\sigma_T} \right)^{2} \right]^{0.5} \]  

(17)

where \( f_b \) is the bottom friction factor, \( \sigma_T \) is the standard deviation of the oscillatory horizontal velocity, and \( |U| = (U^2 + V^2)^{0.5} \).

Since it acts as the “link” between WAVE_L and SEDTR in the framework of the proposed integrated model, it is important for WICIR to be able to reproduce a number of processes that are essential for the realistic description of sediment transport.

Regarding the surf zone, it should be noted that the existence of the undertow (i.e., the current directed offshore) cannot be directly predicted by depth-averaged models; nonetheless, its representation is essential in the aforementioned context. In WICIR, quasi-3D effects are introduced by adopting the analytical expression for the vertical distribution of the cross-shore flow below the wave trough level proposed by Stive and Wind [35], expressed as:

\[ v_u = \frac{1}{2} \left[ (\xi - 1)^2 - \frac{1}{3} \right] \frac{h - \xi_t}{\rho \nu_T} \frac{dR}{dy} + \left( \xi - \frac{1}{2} \right) \frac{h - \xi_t}{\rho \nu_T} - \frac{M \cos \Theta}{h - \xi_t} \]  

(18)

where \( v_u \) is the undertow velocity in the direction normal to the shore, \( \xi = z/(h - \xi_t) \), \( h = d + \xi (\xi \text{ being the mean water elevation}), \xi_t \) is the wave trough level, \( dR/dy = 0.14 \rho g dh/ dy, \tau_t \) is the shear stress at the wave trough level, \( M \) is the wave mass flux above trough level (including surface roller effects), \( \Theta \) is the direction of wave propagation (\( \Theta = \arctan[(<Q_{w2}^2>/<P_{w2}^2>)^{1/2}] \)), and \( \nu_T \) is the eddy viscosity coefficient according to De Vriend and Stive [36]:

\[ \nu_T = 0.025h \left( \frac{D}{\rho} \right)^{1/3} \]  

(19)

Regarding the swash zone, an essential process for shoreline evolution is longshore sediment transport. For obliquely incident waves, the trajectory of the bore-front follows a parabolic movement in the swash, in the direction of the net longshore flow per wave period. The mean longshore transport velocity \( V_R \) at the shoreline is determined according to Baba and Camenen [37] as:

\[ V_R = \sqrt{2gR \sin \Theta} \]  

(20)

where \( R \) is the runup height (\( R = 1.6 H_0 \xi_0 \), where \( H_0 \) is the deep water wave height and \( \xi_0 \) is the Iribarren number) and \( \Theta \) is the wave direction near the rundown point at depth \( d = R/4 \). The longshore velocity \( V_R \) is presumed constant within the swash zone, the width of which is considered as extending from \( d = R/4 \) (i.e., the rundown point) to \( d = -R \). The above velocity is indirectly introduced in the
model by increasing the radiation stresses in the swash zone, based on the rationale described in the following. Longshore velocity can be expressed analytically by:

$$\bar{U} = \frac{2.7}{2} \sqrt{gd_b} \sin a_b \cos a_b$$

(21)

where $\gamma$ is the breaking index, and $d_b$ and $a_b$ are the water depth and incident wave angle at the breaking point, respectively. Assuming a linear variation of $\bar{U}$, the velocity at the shoreline can be approximated as:

$$\bar{U} = 2.7 \gamma \sqrt{d_b} \sin a_b \frac{d_s}{a_b}$$

(22)

where $d_s$ is the water depth at the shoreline. A comparison of Equation (20) to Equation (22) shows that the square of the ratio does not deviate significantly from an empirical factor, $a_s$, expressed as:

$$a_s = 16 \sqrt{\gamma^{1.8} (H_o/L_o)^{0.2}}$$

(23)

where $\gamma_0$ is the Irribarren number, and $H_0$ and $L_0$ are the wave height and wavelength, respectively, for deep water conditions. Accordingly, the aforementioned increase in radiation stresses in the swash zone is achieved by multiplying them by the factor $a_s$.

Finally, regarding flooding due to wave setup, in WICIR, this process is simulated using the “dry bed” boundary condition which, according to Militello et al. [38], can be written as the following set of pairs of conditions for any given grid point $(i,j)$:

- if $(d + \zeta)_{ij} > h_{cr}$ and $(d + \zeta)_{i-1,j} \leq h_{cr}$ and $U_{ij} > 0 \rightarrow U_{ij} = 0$
- if $(d + \zeta)_{ij} > h_{cr}$ and $(d + \zeta)_{i-1,j} \leq h_{cr}$ and $V_{ij} > 0 \rightarrow V_{ij} = 0$
- if $(d + \zeta)_{ij} \leq h_{cr}$ and $(d + \zeta)_{i-1,j} \leq h_{cr} \rightarrow U_{ij} = 0$
- if $(d + \zeta)_{ij} \leq h_{cr}$ and $(d + \zeta)_{i-1,j} \leq h_{cr} \rightarrow V_{ij} = 0$
- if $(d + \zeta)_{ij} \leq h_{cr}$ and $(d + \zeta)_{i-1,j} > h_{cr} \rightarrow U_{ij} = 0$
- if $(d + \zeta)_{ij} \leq h_{cr}$ and $(d + \zeta)_{i-1,j} > h_{cr} \rightarrow V_{ij} = 0$

where $\zeta$ is the mean water surface elevation and $h_{cr}$ is a terminal depth below which drying is assumed to occur (e.g., in WICIR this depth is set to $h_{cr} = 0.001$ m).

The numerical solution in WICIR is also (as in WAVE_L) based on the explicit second order finite difference staggered scheme using a mid-time method [24].

2.3. Sediment Transport Module—SEDTR

The mode of sediment movement on the coast is usually divided into bed load, suspended load, and sheet flow transport. Different model concepts are being presently used for the prediction of each one, which range from empirical transport formulae to more sophisticated bottom boundary layer models. In the present work, bed load transport ($q_b$) is estimated with a quasi-steady, semi-empirical formulation, developed by Camenen, and Larson [19,20] for an oscillatory flow combined with a superimposed current under an arbitrary angle:

$$\Phi_b = \begin{cases} \frac{q_{bw}}{\sqrt{(s-1)g d_{50}}} = a_w \sqrt{\theta_{cw,net} \theta_{cw,m}} \exp\left(-b \frac{d_{50}}{h_{cr}}\right) \\ \frac{q_{bu}}{\sqrt{(s-1)g d_{50}}} = a_u \sqrt{\theta_{cw,m} \theta_{cw,m}} \exp\left(-b \frac{d_{50}}{h_{cr}}\right) \end{cases}$$

(24)

where the subscripts $w$ and $n$ correspond, respectively, to the wave direction and the direction normal to the wave direction; $s = \rho_s/\rho$ is the relative density between the sediment ($\rho_s$) and water ($\rho$); $g$ is the acceleration due to gravity; $d_{50}$ is the median grain size; $a_w$, $a_u$, and $b$ are empirical coefficients; $\theta_{cw,m}$ and $\theta_{cw,m}$ are the mean and maximum Shields parameters due to the wave-current interaction,
respectively; \( \theta_{cw} \) is the current-related Shields parameter in the direction normal to the wave direction, and \( \theta_{cr} \) is the critical Shields parameter for the inception of transport. The net Shields parameter \( \theta_{cw,net} \) in Equation (24) is given by:

\[
\theta_{cw,net} = \left( 1 - a_{pl,b} \right) (1 + a_a) \theta_{cw,on} - \left( 1 + a_{pl,b} \right) (1 - a_a) \theta_{cw,off}
\]

where \( \theta_{cw,on} \) and \( \theta_{cw,off} \) are the mean values of the instantaneous Shields parameter over the two half “periods” \( T_{wc} \) (crest-onshore) and \( T_{wt} \) (trough-offshore), \( a_{pl,b} \) is a coefficient for the phase-lag effects [19], and \( a_a \) is a coefficient for the acceleration effects [39]. The Shields parameter \( \theta_{cw,j} \) is defined by:

\[
\theta_{cw,j} = \frac{1}{2} f_{cw} U_{cw,j}^2 / \left( (s - 1) g d_{50} \right)
\]

with \( U_{cw} \) being the wave and current velocity and \( f_{cw} \) the friction coefficient taking into account the wave and current interaction, while the subscript \( j \) should be replaced either by onshore or offshore. In the above formulation (since linear wave theory cannot be used), the estimation of nonlinear time-varying near-bottom wave velocities is also needed. For the incorporation of nonlinear velocity characteristics (i.e., skewness and asymmetry) in SEDTR, the parameterisation proposed by Isobe and Horikawa [40] is adopted.

The incorporation of the suspended sediment transport rate in SEDTR is done by solving the depth-integrated transport equation for suspended sediment [41,42]:

\[
\frac{\partial (hC)}{\partial t} + \frac{\partial (hCU)}{\partial x} + \frac{\partial (hCV)}{\partial x} = c_{R} w_{s} - w_{s} \beta_{d} \frac{C}{\beta_{d}}
\]

where \( h \) is the total mean depth, \( C \) is the depth-averaged volumetric sediment concentration, \( c_{R} \) is the reference concentration at the bottom [19], \( w_{s} \) is the sediment fall velocity, and \( \beta_{d} \) is a coefficient calculated based on [20] by:

\[
\beta_{d} = \frac{\epsilon}{w_{s}} \left[ 1 - \exp \left( - \frac{w_{s} h}{\epsilon} \right) \right]
\]

with \( \epsilon \) being the sediment diffusivity (related to the eddy viscosity coefficient), estimated by [36]:

\[
\epsilon = 0.025 \ h \left( D / \rho \right)^{1/3}
\]

Cross-shore sediment transport in the swash zone in SEDTR is calculated according to [43], while for the longshore sediment transport, only the increased mean velocity is taken into account, as described in Section 2.2.

These sediment transport rates are then used for the simulation of the coastal bathymetry changes by the module SEDTR. The methodology adopted for the series of model applications can be encoded into the steps described in the following. First, the initial bathymetry is inserted into the wave and wave-induced circulation modules (WAVE_L and WICIR, respectively) in order to estimate the wave and current fields. These fields are afterwards used by the sediment transport module SEDTR to calculate the sediment transport rates. Finally, bathymetry is updated by SEDTR solving the equation of the conservation of sediment transport (for the previously calculated transport rates; [44]). The procedure is repeated for a user-specified time period or until a state of morphologic equilibrium is reached. The aforementioned repetitions take place after bottom change in the order of 10–15% is observed in the field, so that the changes in wave and wave-induced current fields calculated by WAVE_L and WICIR for the updated bathymetry are significant.
3. Model Applications

3.1. Comparison with Experimental Data

The integrated numerical model was set-up and applied in order to reproduce the small-scale laboratory experiments of: (a) Ming and Chiew [21], who studied the shoreline changes caused by the presence of a detached breakwater under the influence of pure wave action; and (b) Badiei et al. [22], who studied the morphological effects of groins on an initially straight beach exposed to oblique irregular waves.

The experiments of Ming and Chiew [21] were conducted in a 10 m long, 5 m wide, and 0.7 m high wave basin. A plunger-type wavemaker was used to generate monochromatic waves and sponge was placed behind the wavemaker in order to minimize wave reflection. The 6 m long beach consisted of uniformly distributed sand with a median grain size of \( d_{50} = 0.25 \) mm. The duration of the tests was approximately 15 h (which was the duration needed for the beach to reach an equilibrium state). Three different cases were reproduced numerically and are presented in the following, for normally incident waves of \( H_0 = 0.05 \) m deep water wave height and \( T = 0.85 \) s wave period. The test cases, presented in Table 1, differed in breakwater length (\( B \)) and breakwater distance from the initial shoreline (\( X \)), in order to cover a wide range of \( B/X \) ratios resulting in both tombolo and salient formation behind the breakwaters.

<table>
<thead>
<tr>
<th>Test</th>
<th>( B ) = Breakwater Length (m)</th>
<th>( X ) = Distance from the Initial Shoreline (m)</th>
<th>( B/X )</th>
<th>Formation of Salient/Tombolo</th>
</tr>
</thead>
<tbody>
<tr>
<td>3</td>
<td>1.5</td>
<td>0.6</td>
<td>2.50</td>
<td>tombolo</td>
</tr>
<tr>
<td>10</td>
<td>1.2</td>
<td>1.2</td>
<td>1.00</td>
<td>salient</td>
</tr>
<tr>
<td>11</td>
<td>1.5</td>
<td>1.2</td>
<td>1.25</td>
<td>tombolo</td>
</tr>
</tbody>
</table>

Badiei et al. [22] employed a series of mobile bed process models (according to [45]) in order to investigate the impact of groins on nearshore morphology under the attack of obliquely incident random waves. Two series of tests were carried out at the Queen’s University Coastal Engineering Laboratory (QUCERL) and the Hydraulic Laboratory of the National Research Council of Canada (NRCC). The physical model regarded an initially plane sloping beach (1:10 slope), composed of \( D_{50} = 0.12 \) mm sand grains. The beach—without the presence of the groins—was exposed to wave action for a duration of 4 h until the formation of a nearly stable bathymetry (clear offshore bar trough/step formation). The installation of the groins followed, and the tests continued thereafter in 2 h cycles. In this work, the case of a single groin was modelled, exposed to waves of \( H_s = 0.08 \) m deep water significant wave height, \( T_p = 1.15 \) s peak period, and \( \theta_0 = 11.6^\circ \) deep water incident wave angle, for a total duration of 12 h after groin installation (Test NT2, NRCC test series).

3.2. Application to Paralia Katerinis Beach (Greece)—Coastal Protection with Submerged Breakwaters

Following its validation for laboratory experiments, the presented integrated model was applied to a real case study of a coastal engineering project regarding the protection of a beach by using detached submerged breakwaters. The study area is located in the Region of Central Macedonia, Greece, at a sandy beach north of the fishing port of Katerini (Figure 1). The area has been facing coastal erosion problems for well over 30 years, which started after the construction of the harbour seen in the left part of Figure 1 (1980–1984). As a result, and due to the prevailing SE winds, the coastal zone south of the fishing port showed strong accretion, while the coast north from the port was eroded, with a shoreline retreat in the order of 20 m; the erosive phenomena stretched over a zone of 800 m north from the port. In order to reverse erosion, a groin field consisting of 13 rubble mount groins was constructed (1990–1997). The project not only failed to further protect the beach—since additional erosion occurred in between the constructed groins—but it also transferred coastal retreat northwards.
Furthermore, the semi-closed basins that formed between the constructed groins (being in the non-tidal Mediterranean Sea) caused significant environmental problems regarding water quality due to the limited renewal rates. In 2010, a new coastal protection project was designed and constructed; the groin field was replaced by a set of three 200 m long submerged breakwaters placed at a distance of approximately 200 m from the coast, and the gaps between the structures were approximately 110 m. The breakwaters were designed to have transmission coefficients in the order of 0.4 ($K_t \approx 0.4$). In addition, a beach nourishment project was also designed and applied to restore the beach to its previous condition. Bathymetry measurement data for the area are available for the period right after the completion of coastal works (beach nourishment and construction of the submerged breakwaters), as well as for three years later [46].

![Location and satellite image of the study area at Paralia Katerinis, Greece (privately processed).](image)

The integrated model was applied to simulate coastal morphodynamics after the realization of the coastal protection project. Since wave data were not available for the area, hindcast data were used. The main incident wind directions are: NE, E, and SE-S. The model was run by applying three representative waves (i.e., three equivalent wave heights on an annual basis; see Table 2). The workflow for the coupled module runs can be summed-up in the following. Starting with the initial bathymetry, WAVE_L, WICIR and SEDTR modules were run in sequence for the characteristics of the first representative wave and taking into account its annual frequency of occurrence for the simulation of morphology evolution. The updated bed morphology was then used to run the integrated model for the second representative wave (in the same way) and the morphology at the end of this second run for the third representative wave. The aforementioned simulation steps were repeated until the total duration of the wave action was reached.

Table 2. Characteristics of the three representative waves used for the Paralia Katerinis beach runs.

<table>
<thead>
<tr>
<th>Wave Direction</th>
<th>Significant Wave Height $H_s$ (m)</th>
<th>Peak Wave Period $T_p$ (s)</th>
<th>Annual Frequency of Occurrence $f$ (%)</th>
</tr>
</thead>
<tbody>
<tr>
<td>SE-S</td>
<td>1.34</td>
<td>5.4</td>
<td>12.70</td>
</tr>
<tr>
<td>E</td>
<td>1.09</td>
<td>4.4</td>
<td>1.12</td>
</tr>
<tr>
<td>NE</td>
<td>0.94</td>
<td>4.6</td>
<td>1.17</td>
</tr>
</tbody>
</table>

4. Results and Discussion

4.1. Comparison with Experimental Data

Figure 2 shows the initial wave-induced current velocity field and the comparison between the computed and measured shoreline evolution data for Test 11 of Ming and Chiew [21]. The presence of the breakwater leads to the formation of two opposing eddies in the area behind it, as currents move

![Figure 2 shows the initial wave-induced current velocity field and the comparison between the computed and measured shoreline evolution data for Test 11 of Ming and Chiew [21].](image)
towards the sheltered area along the foreshore from both sides of the structure. A secondary cause of the observed circulation pattern is the mean sea level gradient between the illuminated and sheltered areas due to diffraction effects, while the representation of swash zone hydrodynamics by the model should also be highlighted (see also Section 2.2). Regarding morphology evolution, the model results are in good agreement with the measured data of Ming and Chiew [21], satisfactorily representing the formation of the tombolo in Test 11 (Figure 2b), as well as the formation of the tombolo and salient in Tests 3 (Figure 3a) and 10 (Figure 3b), respectively.

Figure 2. (a) Initial wave-induced current velocity field (contours represent the initial bathymetry) and (b) comparison between the computed and measured shoreline evolution (contours represent the final computed bathymetry) for Test 11 of Ming and Chiew [21].
Figure 3. Comparison between the computed and measured shorelines evolution (contours represent the final computed bathymetry) for: (a) Test 3 and (b) Test 10 of Ming and Chiew [21].

Figure 4 shows the initial breaking wave-induced current velocity field and the comparison between the computed and measured shoreline evolution for the single groin test (Test NT2) of Baidei et al. [22]. The sediment accretion updrift of the groin-type structure results in an advance of the shore, while the lack of sediments at the lee of the groin leads to a retreat of the shore; shoreline evolution is well-reproduced, with model results and measurement lines practically overlapping.

In general, the integrated model results in smooth but consistent (considering also test runs that are not presented in this work) bathymetries behind detached breakwaters and in the vicinity of groin-type structures, while it appears to be smoothing in a close-to-natural way eventual local shoreline irregularities which, on the other hand, are present in the experimental results.
4.2. Application to Paralia Katerinis Beach (Greece)—Coastal Protection with Submerged Breakwaters

Figure 5 shows the breaking wave-induced current field for the prevailing SE-S waves. The submerged breakwaters allow some wave transmission and overtopping that cause an additional supply of water behind the structures, which is taken into account by the circulation module and affects the current field. This net transport of water into the lee zone causes a water level rise and is balanced mainly by outgoing currents at the heads of the structures. Consequently, the main flow pattern is characterized by an onshore flow over the submerged breakwaters, an offshore flow at the gaps between them (eroding rip currents), and nearshore eddies similar to those formed in the case of emerged breakwaters (although with lower relative intensity). The first two flow patterns do not exist in the latter case, while the third flow pattern is not extended up to the structures, where onshore flow leads to the formation of—more intense—opposite direction eddies.

Figure 6 shows the initial bed morphology of Paralia Katerinis beach and the comparison between the computed and measured bed morphology evolution, along with the satellite image of the same area (2016). The workflow for the coupled module runs is described in Section 3.2. The model results are in close agreement with the measurements (respective lines are practically overlapping), with the integrated model succeeding in reproducing all morphological patterns behind the breakwaters and up to the shoreline, under the presence of both permanent structures (emerged groins) and the realized beach nourishment project there.
Figure 4. (a) Initial breaking wave-induced current velocity field (contours represent the initial bathymetry) and (b) comparison between the computed and measured shoreline evolution (contours represent the final computed bathymetry) for Test NT2 of Baidei et al. [22].

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Figure 5. Breaking wave-induced current field for the prevailing SE waves at Paralia Katerinis beach (Greece) after the construction of the detached submerged breakwaters.

Figure 6 shows the initial bed morphology of Paralia Katerinis beach and the comparison between the computed and measured bed morphology evolution, along with the satellite image of the same area (2016). The workflow for the coupled module runs is described in Section 3.2. The model results are in close agreement with the measurements (respective lines are practically overlapping), with the integrated model succeeding in reproducing all morphological patterns behind the breakwaters and up to the shoreline, under the presence of both permanent structures (emerged groins) and the realized beach nourishment project there.

Figure 6. (a) Initial bed morphology of Paralia Katerinis beach, (b) comparison between the computed and measured bed morphology evolution, and (c) satellite image of the same area ([47]; privately processed).

4.3. General Discussion

Elaborating further on the integrated model’s performance in simulating the morphological effects of the presence of coastal protection structures, particular insights can be drawn by the discussion in the following.

Model setup was based on the successful model calibration for the experimental data presented in this work (see Sections 3.1 and 4.1), along with specific modelling choices based on extensive experience in morphological modelling for both research and engineering applications. The calibration process of the presented model mainly refers to: (a) swash zone incorporation, i.e., the...
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The satisfactory agreement between model predictions and experimental/field data should be mainly attributed to the above choices, which are deemed to distinguish the presented model from relevant work. The novelty of incorporating swash zone hydrodynamics using a linear wave model and a nonlinear wave-induced circulation model particularly allows for the simulation of swash zone morphodynamics which play an essential role in nearshore morphology evolution. The above, keeping in mind that the objective of this work was to present an integrated model that could be adapted to a wide range of design layouts and perform satisfactorily for an equally wide range of field conditions, while maintaining the computational effort needed at reasonable levels, an aspect that usually limits the applicability of more complex models in engineering applications.

5. Conclusions

This work presents an integrated coastal engineering numerical model that simulates linear wave propagation, wave-induced circulation, sediment transport, and bed morphology evolution. The model consists of three main modules: The nearshore wave transformation model WAVE_L, properly adapted for the simulation of compound linear wave fields near coastal structures; the wave-induced circulation module WICIR, which includes a novel approach for the representation of swash zone hydrodynamics; and the sediment transport and bed morphology evolution module SEDTR.

The model is tested against experimental data to study the effect of representative coastal protection structures, such as detached breakwaters (data from [21]) and groins (data from [22]). Given the good agreement between model results and laboratory measurements, the model was also successfully applied to a real case study of a coastal engineering project in North Greece (combination of submerged breakwaters and beach nourishment). The model is deemed to constitute a suitable tool for the design and evaluation of the morphological influence of harbour and coastal protection works, being able to deliver results in a fast and seamless way at all times for a wide range of design layouts.

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References and Note


